

## A variational principle in non-smooth Mechanics

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Let  $Y$  be a Banach space with dual  $Y^*$ ,  $\psi : Y \rightarrow (-\infty, \infty]$  be convex, proper, and lower semicontinuous dissipation potential,  $\phi \in C^1(Y; \mathbb{R})$  be an energy, and the load  $\ell \in L^\infty(0, T; Y^*)$  and the initial state  $y_0 \in Y$  be given. A variety of non-smooth (thermo-)mechanical evolution models can be represented by the relation

$$(1) \quad \partial\psi(\dot{y}) + D\phi(y) \ni \ell \quad \text{a.e. in } (0, T), \quad y(0) = y_0,$$

where  $\partial$  stands for the subdifferential in the sense of Convex Analysis and  $t \mapsto y(t)$  is meant to be in  $W^{1,1}(0, T; Y)$ . Here  $\partial\psi$  represents the system of dissipative forces whereas  $D\phi - \ell$  stands for the conservative forces instead. Relation (1) appears in connection with many different applicative situations among which plasticity, visco-elasticity, friction [6], heat conduction, and phase change.

Fenchel's inequality  $\psi(u) + \psi^*(v) \geq \langle v, u \rangle$  holds for all  $u \in Y$  and  $v \in Y^*$  where  $\langle \cdot, \cdot \rangle$  stands for the duality pairing between  $Y^*$  and  $Y$  and  $\psi^*$  is the conjugate of  $\psi$ . Moreover, the latter inequality reduces to an equality iff  $v \in \partial\psi(u)$ . Hence, by defining the *Lagrangian*  $L : (0, T) \times Y \times Y \rightarrow \mathbb{R}$  as

$$L(t, y, p) := \psi(p) + \psi^*(\ell(t) - D\phi(y)) - \langle \ell(t) - D\phi(y), p \rangle$$

for almost every  $t \in (0, T)$ , one readily checks that  $L(t, y, p) \geq 0$  and

$$L(t, y, p) = 0 \quad \text{iff} \quad \partial\phi(p) \ni \ell(t) - D\phi(y).$$

Let now the functional  $F : W^{1,1}(0, T; Y) \rightarrow [0, \infty]$  be defined as

$$F(y) := \int_0^T L(t, y(t), \dot{y}(t)) dt + \phi(y(0) - y_0).$$

Then, minimizers of  $F$  and solutions to (1) coincide. Indeed, as  $y$  solves (1), we have that  $L(t, y(t), \dot{y}(t)) = 0$  for almost every  $t \in (0, T)$  and  $y(0) = y_0$ . That is  $F(y) = 0$ . On the contrary, if  $F(y) = 0$ , then  $L(t, y(t), \dot{y}(t)) = 0$  for almost every  $t \in (0, T)$  and  $\phi(y(0) - y_0) = 0$  and (1) follows. Hence, we have the following.

**Theorem 1.**  *$y$  solves (1) iff  $F(y) = 0 = \min F$ .*

The characterization of solutions of differential problems driven by convex potentials as minimizers of functionals via the Fenchel approach is rather classical and has to be traced back to Brezis & Ekeland [1, 2]. In the referred papers, the Authors provide some similar variational characterization for gradient flows of convex functionals in Hilbert spaces. I have recently extended these ideas to the case of doubly nonlinear equations [7].

In order to extract some information from the variational characterization of Theorem 1, it is convenient to focus on some more specific situation. I shall in particular consider the case of linearized hardening elasto-plasticity [4] by additionally requiring

$$\psi \text{ positively 1-homogeneous and } \phi \text{ quadratic and coercive on } D(\psi),$$

namely the domain of  $\psi$ . In this specific situation  $D\phi$  is linear and Theorem 1 is providing a useful tool in order to deal with limiting procedures. Indeed, since one reinterprets the equation as a minimum problem, it is clear that the natural concept to be considered is that of  $\Gamma$ -convergence [3]. More precisely, as Theorem 1 directly quantifies the value of the minimum to be 0, what is actually needed are so-called  $\Gamma$ -lim inf inequalities only.

The variational approach by Theorem 1 provides the possibility of recovering convergence in a variety of approximated situations.

**Space approximation.** Conformal finite elements and data approximations can be considered. In particular, one can easily handle by means of this variational approach the case where  $Y$  is replaced by a nested sequence of finite dimensional subspaces  $Y_h$  such that  $\cup_h Y_h$  is dense in  $Y$ ,  $\phi_h$  is the restriction of  $\phi$  to  $Y_h$ , and  $\psi_h \rightarrow \psi$  in the sense of Mosco in  $Y$  [3].

**Time discretization.** One can provide a discrete analogue to Theorem 1. Namely, it is possible to variationally characterize solutions of the so-called  $\theta$ -scheme

$$y^0 = y_0, \quad \psi(y^i - y^{i-1}) + A(\theta y^i + (1 - \theta)y^{i-1}) \ni \ell_\theta^i \quad \text{for } i = 1, \dots, N,$$

where  $\ell_\theta^i = \ell(\theta t^i + (1 - \theta)t^{i-1})$  for  $0 = t^0 < t^1 < \dots < t^N = T$ , as minimizers of a suitable discrete functional tailored on  $F$ . This approach can be exploited in order to prove the unconditional stability of the scheme in  $W^{1,\infty}(0, T; Y)$  and its convergence with respect to the corresponding weak-star topology as the diameter of the partition goes to 0. Finally, the discrete functional itself can be used in order to control the discretization error.

**Fully discrete schemes.** The above-mentioned techniques for space and time discretization can be combined in order to provide the weak-star convergence in  $W^{1,\infty}(0, T; Y)$  of some fully discretized solutions as well. This result has to be compared with the former analysis in [5], where a stronger convergence is proved by means of a somehow more complicated argument.

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